## Measures on Operational Logics\*

Sławomir Bugajski \*\*

Institute of Physics, Silesian University, Katowice, Poland

Z. Naturforsch. 34a, 785-786 (1979); received March 19, 1979

It is demonstrated, that any measure on effects is generated by a state, whereas a similar statement for operations is not true.

The operational approach to quantum mechanics [1, 2, 3] introduces at least two sets which to some extent could be regarded as generalizations of the "quantum logic": the set of effects (also simple observables, tests), and the set of operations. Both sets carry natural structures making it possible to define generalized measures on them. Moreover, any state induces such measures, hence the question arises: is any measure on effects (resp. operations) generated by some state? This problem, formulated by R. S. Ingarden ([4], p. 115) is solved below.

We start with three postulates [5] which underlie the operational approach to quantum mechanics.

Postulate 1. The set of states of a physical system is represented by the set of elements of a closed generating cone S for a complete base norm space  $(V, S_1)$ .

Postulate 2. The set  $L_{\text{ef}}$  of effects is represented by the unit interval  $\langle o, e \rangle$  in the dual space  $(V^*, e)$  of  $(V, S_1)$ .

Postulate 3. The set  $L_{op}$  of operations on the physical system is represented by the set of positive elements in the unit ball of  $\mathcal{L}(V)$ .

We define bounded measures on  $L_{\mathrm{ef}}$  [4] as mappings  $\mu \colon L_{\mathrm{ef}} \mapsto \mathbb{R}$  such that:

- (i)  $\mu(a) \ge 0$  for any  $a \in L_{ef}$ ,
- (ii)  $\mu(0) = 0$ ,  $\mu(e) < \infty$ ,

(iii) 
$$\sum_{k=1}^{\infty} a_k \in L_{\text{ef}} \mid \Rightarrow \mu \left( \sum_{k=1}^{\infty} a_k \right) = \sum_{k=1}^{\infty} \mu \left( a_k \right)$$

for any sequence of effects with the sum converging in the weak\* topology.

0340-4811 / 79 / 0600-0785 \$ 01.00 / 0

Any state (i.e. any element of the cone S) generates a weak\* continuous positive linear functional on  $V^*$ , which when restricted to  $L_{\rm ef}$  defines a measure in the above sense. It is evident, however, that it suffices to take a sequentially weak\* continuous positive functional on  $V^*$  to obtain a measure on  $L_{\rm ef}$ . So our problem reduces to the question: are there positive linear sequentially weak\* continuous functionals on  $V^*$  which are not weak\* continuous?

Let  $\alpha$  be such a functional, then  $\alpha^{-1}(0)$  is a sequentially weak\* closed, but weak\* dense, hypersubspace of  $V^*$ . As the positive cone  $S^*$  of  $V^*$  has non-empty weak\* interior, then  $\alpha^{-1}(0) \cap S^*$  cannot be empty. It suffices for the conclusion: any positive linear sequentially weak\* continuous functional on  $V^*$  is also weak\* continuous, hence any measure on  $L_{\rm ef}$  is generated by a state.

One defines measures on  $L_{op}$  [3, 4] as mappings  $\mu: L_{op} \rightarrow \mathbb{R}$  such that:

- (i)  $\mu(\Phi) \ge 0$  for any  $\Phi \in L_{op}$ ,
- (ii)  $\mu(\Omega) = 0$ ,  $\mu(I) < \infty$ ,

(iii) 
$$\sum_{k=1}^{\infty} \Phi_k \in L_{\text{op}} \mid \Rightarrow \mu \left( \sum_{k=1}^{\infty} \Phi_k \right) = \sum_{k=1}^{\infty} \mu \left( \Phi_k \right)$$

for any sequence of operations with the sum converging in the strong operator topology, where  $\Omega \alpha = \omega$  (the zero-vector of V) for any  $\alpha \in V$ ,  $I\alpha = \alpha$  for any  $\alpha \in V$ .

Any linear positive sequentially strong continuous functional on  $\mathscr{L}(V)$  generates measure on  $L_{\mathrm{op}}$ , and conversely. States generate standard weak continuous functionals  $\Phi \to e(\Phi \alpha)$ ,  $\alpha \in S$ . Observe that these functionals do not separate elements of  $\mathscr{L}(V)$ . As the family  $\{a(\cdot \alpha) \mid a \in V^*, \alpha \in S, a \geq 0\}$  of positive weak continuous functionals on  $\mathscr{L}(V)$  does separate elements of  $\mathscr{L}(V)$ , we conclude that the set of measures on  $L_{\mathrm{op}}$  is far larger than the set of state generated ones.

In the standard von Neumann model of quantum mechanics we start with a von Neumann algebra  $\mathscr{B}$  on a Hilbert space. The space  $(V, S_1)$  is then realized as the space of all self-adjoint ultra-weak continuous linear functionals on  $\mathscr{B}$ , with  $S_1$  — the convex set of normalized normal positive functionals on  $\mathscr{B}$ . The strictly positive linear functional e on V, corresponding to the base  $S_1$ , is just the trace.  $V^*$  can be identified with the set of self-adjoint elements of  $\mathscr{B}$ . It is rather obvious now that in this model any measure on  $L_{\rm ef}$  is generated



Dieses Werk wurde im Jahr 2013 vom Verlag Zeitschrift für Naturforschung in Zusammenarbeit mit der Max-Planck-Gesellschaft zur Förderung der Wissenschaften e.V. digitalisiert und unter folgender Lizenz veröffentlicht: Creative Commons Namensnennung-Keine Bearbeitung 3.0 Deutschland

This work has been digitalized and published in 2013 by Verlag Zeitschrift für Naturforschung in cooperation with the Max Planck Society for the Advancement of Science under a Creative Commons Attribution-NoDerivs 3.0 Germany License.

<sup>\*</sup> Supported by Polish Ministry of Science, Higher Education and Technology Project MR I 7.

<sup>\*\*</sup> Reprint requests to S. Bugajski, Instytut Fizyki U. Sl, ul. Uniwersytecka 4, 40-007 Katowice, Poland.

by a state (is regular, in terms of Ingarden). This fact could be regarded as a generalization of the Gleason theorem. Our general conclusion concerning measures on operations holds also in this case.

It would be of interest to study states on a special class of operations, called filters (or strongly repeatable operations), as well on the corresponding effects, called decision effects. This is not an easy task, however, as there is no agreement about the precise notion of filter (compare definitions in [2], p. 200, [1], p. 254, and a discussion in [6], where further references can be found).

- [1] E. B. Davies and J. T. Lewis, Commun. Math. Phys. 17, 239 (1970).
- [2] A. Hartkämper and H. Neumann (eds.), Foundations of Quantum Mechanics and Ordered Linear Spaces. Advanced Study Institute Marburg 1973, Springer-Verlag, Berlin 1974.
- M. D. Srinivas, J. Math. Phys. 16, 1672 (1975).R. S. Ingarden, Information Theory and Thermodynamics Part I Ch. II, Preprint of the Copernicus University No. 275, Torun 1976.
- [5] C. M. Edwards, Commun. Math. Phys. 16, 207 (1970).
  [6] W. M. Cornette and S. P. Gudder, J. Math. Phys. 15, 842 (1974).

## Berichtigung

Chr. Bräuchle, H. Kabza, and J. Voitländer, Optical and ODMR Investigations of the Lowest Excited Triplet State of Dinaphtho-(2'.3':1.2); (2".3":6.7)-pyrene, Z. Naturforsch. 34a, 6 [1979].

In dieser Arbeit müssen die Fußnoten [6] und [7] wie folgt lauten:

- [6] Chr. Bräuchle, H. Kabza, J. Voitländer, and E. Clar, Chem. Phys. 32, 63 (1978).
- [7] E. Clar, W. Schmidt, and T. Zauner, to be published.